

# Energy and entropy in General Relativity

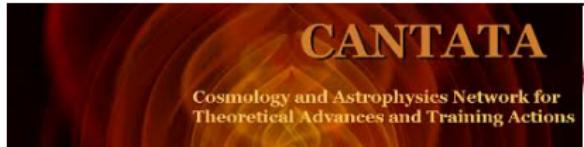
## - or, the Canonical Frame in Q-Gravity

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# Displacement current of the metric

R. P. Feynman, *Feynman lectures on gravitation* (1962)

“Thus, gravity is that field which corresponds to a gauge invariance with respect to displacement transformations.”

- Recall the original\* EMT of the gravitational field:

$$t_{E\nu}^{\mu} = \frac{\partial L_E}{\partial g_{\alpha\beta,\nu}} g_{\alpha\beta,\nu} - \delta_{\nu}^{\mu} L_E .$$

- It is also the canonical<sup>†</sup> displacement current.

\* A. Einstein, *Hamiltonsches Prinzip und allgemeine Relativitätstheorie* (1916)

† E. Noether, *Invariante Variationsprobleme* (1918)

# Pseudo-tensors and complexes

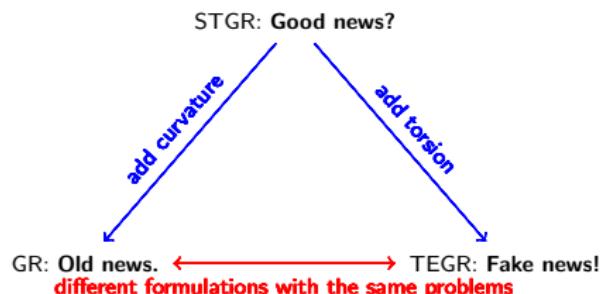
A. Einstein, *Hamiltonsches Prinzip und allgemeine Relativitätstheorie* (1916):

"Durch Ausführung der partiellen Integration"  $\mathcal{R} \rightarrow$

$$L_E = g^{\mu\nu} \left( \left\{ \frac{\alpha}{\sigma\mu} \right\} \left\{ \frac{\sigma}{\nu\alpha} \right\} - \left\{ \frac{\alpha}{\sigma\alpha} \right\} \left\{ \frac{\sigma}{\mu\nu} \right\} \right)$$

- The use of  $\tau_E^\mu{}_\nu = t_E^\mu{}_\nu + T^\mu{}_\nu$  can be criticised<sup>‡</sup>.
- Many non-canonical alternatives have been proposed<sup>§</sup>.

- In particular: metric teleparallelism does not remove, but replaces the problem.



<sup>‡</sup> Levi-Civita, Schrödinger, Weyl, Eddington, Pauli, etc

<sup>§</sup> Tolman, Landau-Lifshitz, Papapetrou-Gubta, Møller, Weinberg, etc

# Affine connection

- **Curvature:**  $R^\alpha_{\beta\mu\nu} = 2\partial_{[\mu}\Gamma^\alpha_{\nu]\beta} + 2\Gamma^\alpha_{[\mu|\lambda|}\Gamma^\lambda_{\nu]\beta}$
- **Torsion:**  $T^\alpha_{\mu\nu} = 2\Gamma^\alpha_{[\mu\nu]}$
- **Non-metricity:**  $Q_{\alpha\mu\nu} = \nabla_\alpha g_{\mu\nu}$

$$\overbrace{\Gamma^\alpha_{\mu\nu}}^{\text{affine connection}} = \underbrace{\left\{ \begin{matrix} \alpha \\ \mu\nu \end{matrix} \right\}}_{\text{Levi-Civita}} + \underbrace{K^\alpha_{\mu\nu}}_{\text{contortion}} + \underbrace{L^\alpha_{\mu\nu}}_{\text{disformation}}$$
$$-\frac{1}{2}g_{\mu\nu,\alpha} + g^{\alpha\lambda}g_{\lambda(\mu,\nu)} - \frac{1}{2}T^\alpha_{\mu\nu} - T_{(\mu\nu)}^\alpha - \frac{1}{2}Q^\alpha_{\mu\nu} - Q_{(\mu\nu)}^\alpha$$

## Geometrical interpretation

- $R_{[ab]} = \frac{1}{2}\partial_a^\alpha\partial_b^\beta R_{[\alpha\beta]\mu\nu}dx^\mu\wedge dx^\nu$ : how things turn around
- $T^a = \frac{1}{2}\epsilon^a_\alpha T^\alpha_{\mu\nu}dx^\mu\wedge dx^\nu$ : how things are displaced
- $R_{(ab)}$ : how things change their size and shape

# Affine curvature

- The curvature scalar can be decomposed accordingly:

$$\underbrace{R}_{\text{affine}} = \underbrace{\mathcal{R}}_{\text{metric}} + \underbrace{\mathbb{S}}_{\text{stuff}} + \underbrace{\mathcal{D}_\mu}_{\text{metric}} \left( \underbrace{Q^\mu}_{Q^{\mu\alpha}{}_\alpha} - \underbrace{\tilde{Q}^\mu}_{Q^{\alpha\mu}{}_\alpha} + 2 \underbrace{T^\mu}_{T^{\alpha\mu}{}_\alpha} \right)$$

- Therefore, for a flat connection with  $R = 0$ , we have

$$\underbrace{\frac{1}{2}M^2 \int d^4x \sqrt{-g} \mathcal{R}}_{\text{Einstein-Hilbert}} \underset{\text{up to a boundary term}}{=} \underbrace{-\frac{1}{2}M^2 \int d^4x \sqrt{-g} \mathbb{S}}_{\text{Teleparallel equivalent}}$$

The post-Riemannian Scalar is the quadratic term

$$\begin{aligned} \mathbb{S} = & \frac{1}{4} T_{\mu\nu\rho} T^{\mu\nu\rho} + \frac{1}{2} T_{\mu\nu\rho} T^{\mu\rho\nu} - T_\mu T^\mu + Q_{\mu\nu\rho} T^{\rho\mu\nu} - Q_\mu T^\mu + \tilde{Q}_\mu T^\mu \\ & + \frac{1}{4} Q_{\mu\nu\rho} Q^{\mu\nu\rho} - \frac{1}{2} Q_{\mu\nu\rho} Q^{\nu\mu\rho} - \frac{1}{4} Q_\mu Q^\mu + \frac{1}{2} Q_\mu \tilde{Q}^\mu \end{aligned}$$

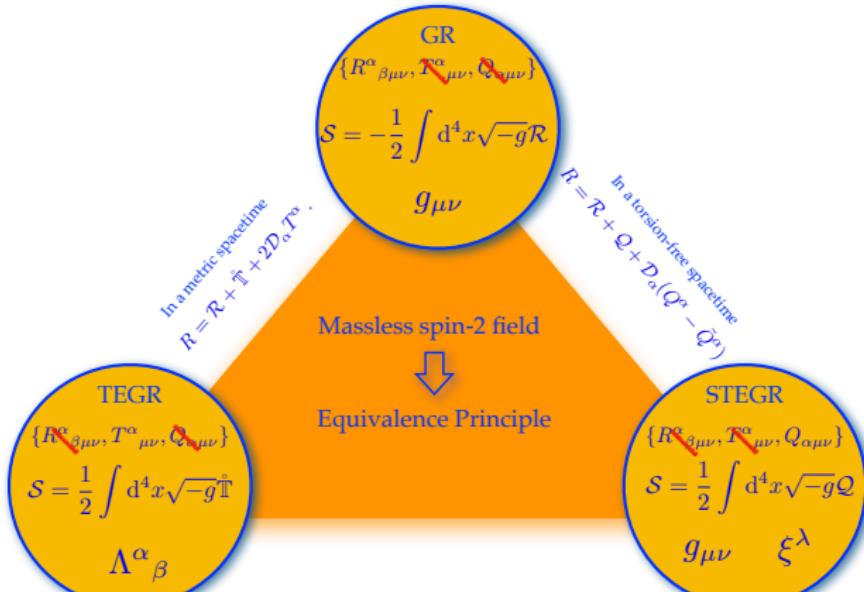
[Jose Beltrán Jiménez, Lavinia Heisenberg, Damianos Iosifidis, Alejandro Jiménez-Cano, TSK: arXiv:1909.09045]



# The geometrical trinity

[Jose Beltrán Jiménez, Lavinia Heisenberg, TSK: Universe 5 no.6, 153 (2019)]

$$\# \text{dof's} = 10 - 2 \times (\text{4 Diffs}) = 2$$



$$\# \text{dof's} = 16 - 2 \times 4 \text{ Diffs} - 6 \text{ (Local Lorentz)} = 2 \quad \# \text{dof's} = 10 + 4 - 4 \text{ Diffs (CG)} - 2 \times 4 \text{ Diffs' (bound.)} = 2$$

# (Re)interpretations

## General Relativity

$$\ddot{x}^\alpha + \overbrace{\left\{ \begin{array}{c} \alpha \\ \mu\nu \end{array} \right\}}^{\text{geometry}} \dot{x}^\mu \dot{x}^\nu = 0.$$

## Teleparallelism

$$\ddot{x}^\alpha + \overbrace{\Gamma^\alpha_{\mu\nu}}^{\text{Weitzenböck}} \dot{x}^\mu \dot{x}^\nu = \overbrace{K^\alpha_{\mu\nu} \dot{x}^\mu \dot{x}^\nu}^{\text{force } F^\alpha}.$$

## Symmetric teleparallelism

$$\ddot{x}^\alpha + \overbrace{\Gamma^\alpha_{\mu\nu}}^{\text{pure gauge}} \dot{x}^\mu \dot{x}^\nu = \overbrace{L^\alpha_{\mu\nu} \dot{x}^\mu \dot{x}^\nu}^{\text{inertia } I^\alpha}.$$



# Translation connection

- The **generic** affine connection  $\Gamma^\alpha_{\mu\nu}$ : 64 functions.
- The **flat**:  $\Gamma^\alpha_{\mu\nu} = \Lambda^\alpha_\beta \partial_\mu (\Lambda^{-1})^\beta_\nu$ : 16 functions.
- The **flat and metric-compatible**:  $\Lambda^{\alpha\beta} = \Lambda^{[\alpha\beta]}$ : 6 functions.
- The **flat and symmetric**:  $\Lambda^\alpha_\beta = \partial_\beta \xi^\alpha$ : 4 functions.

Jose Beltrán Jiménez, Lavinia Heisenberg, TSK: *Coincident General Relativity* (2017):

The affine connection of “**purified gravity**” is a pure translation; it has no geometry in the gauge-invariant sense. The connection vanishes in the “**coincident gauge**”,  $\mathring{\nabla}_\mu = \partial_\mu$ .

# Coincident General Relativity

[J. Beltrán Jiménez, L. Heisenberg, TSK: PRD98, no.4 (2018)]

- There is a **unique quadratic form** which is doubly translation-invariant:  $Q = -Q_{\alpha}^{\mu\nu}P_{\mu\nu}^{\alpha}$ .

$$P^{\alpha}_{\mu\nu} = -\frac{1}{4}Q_{\alpha}^{\mu\nu} + \frac{1}{2}Q_{(\mu}^{\alpha}\nu) + \frac{1}{4}(Q^{\alpha} - \tilde{Q}^{\alpha})g_{\mu\nu} - \frac{1}{4}\delta_{(\mu}^{\alpha}Q_{\nu)}, \text{ where } Q_{\alpha} = Q_{\alpha}^{\mu}\mu, \tilde{Q}^{\alpha} = Q_{\mu}^{\alpha\mu}.$$

- At linear order, the connection decouples from the action.
- In the coincident gauge,  $Q$  reduces to  $L_E$ ,  $\mathring{Q} = -L_E$ .

Recall, "durch Ausführung der partiellen Integration [...]" (1916)

$$\mathcal{R} \xrightarrow{\text{part. int.}} L_E = g^{\mu\nu} \left( \left\{ \begin{matrix} \alpha \\ \sigma\mu \end{matrix} \right\} \left\{ \begin{matrix} \sigma \\ \nu\alpha \end{matrix} \right\} - \left\{ \begin{matrix} \alpha \\ \sigma\alpha \end{matrix} \right\} \left\{ \begin{matrix} \sigma \\ \mu\nu \end{matrix} \right\} \right) = -\mathring{Q} \xleftarrow{\text{gauge fix}} -Q$$

# The coupling matters in a gauge theory.

The electromagnetic field:

$$F = \partial_{[\mu} A_{\nu]} dx^\mu \wedge dx^\nu \quad \xrightarrow{\text{M.C.}} \quad F - \frac{1}{2} T^\alpha{}_{\mu\nu} A_\alpha dx^\mu \wedge dx^\nu .$$

Y. So, J. Nester, *On the source coupling and the teleparallel equivalent to GR* (2006):

“Can one decide which is the real ‘physical’ geometry? Invoking the minimal coupling principle may give a unique answer.”

J. Itin, Yu. Obukhov, J. Boos, F. Hehl, *Premetric teleparallel theory of gravity and its local and linear constitutive law* (2018):

“Within the framework of TG as a translational gauge theory, the coupling to matter is achieved via the minimal coupling procedure, strictly in the sense of a bona fide gauge theory. Dispensing with the minimal coupling principle [...] is against the spirit of gauge theory.”

# Coupling matter: works in PG, doesn't in TG.

The hermitean Dirac action:

$$S_D = -\frac{1}{2} \int d^4x \sqrt{-g} \left[ (i\bar{\psi}\gamma^\mu D_\mu \psi) + (i\bar{\psi}\gamma^\mu D_\mu \psi)^\dagger - 2m\bar{\psi}\psi \right].$$

- Tangent space:  $D_\mu e^a = 0 \Rightarrow \Lambda^a{}_{\mu b} = e^a{}_\nu (\nabla_\mu \theta_b{}^\nu)$ .
- Spin connection:  $\Gamma = \frac{1}{4} A_{ab} \gamma^a \gamma^b + \frac{1}{8} Z = \frac{1}{4} A_{ab} \gamma^{[a} \gamma^{b]} + \frac{1}{8} Q + \frac{1}{8} Z$ .
- Hermiteanised:  $\gamma \cdot \Gamma^H = -\frac{i}{4} \epsilon^{abcd} \text{Re}(A_{ab}) \cdot \theta_c \gamma_d \gamma^5 + \frac{i}{2} \text{Im}(A^{[ab]}) \cdot \theta_a \gamma_b - \frac{i}{8} \text{Im}(Q + Z) \cdot \gamma$ .
- Decomposed:  $i\Gamma_\mu^H = \frac{1}{4} \text{Re}(\tilde{\omega}_\mu + \tilde{T}_\mu) \gamma^5 - \frac{1}{2} \text{Im}(\omega_\mu - T_\mu + \frac{1}{2} \tilde{Q}_\mu) + \frac{1}{8} \text{Im}(3Q_\mu + Z_\mu)$ .

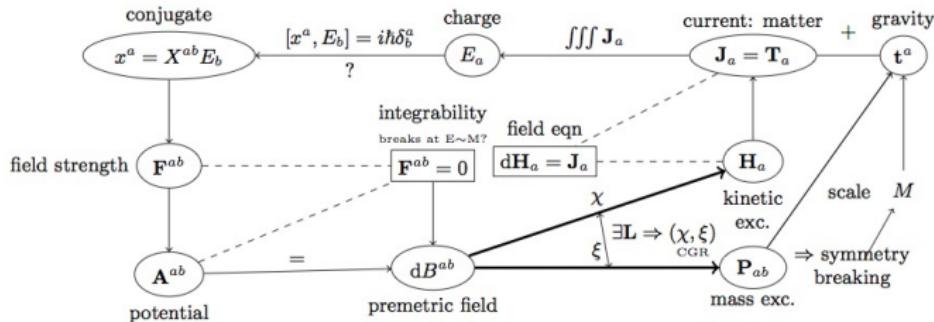
The hermitean Dirac equation:

$$(i\gamma^\mu D_\mu^H - m) \sqrt{-g} \psi = 0.$$

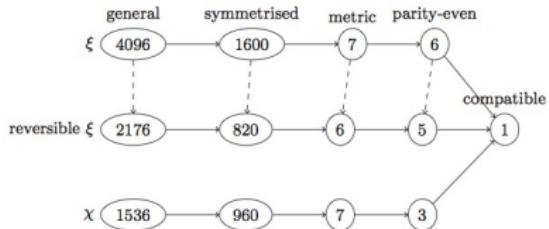
The point is that though  $\Gamma = \{ \} + L \stackrel{?}{=} 0$ , spinors see only the  $\{ \}$ .



# Premetric Q-gravity [TSK, Manuel Hohmann, Luca Marzola: arXiv:1909.10415].



- **Fundamental equations**
  - conservation:  $d\mathbf{H}_a = \mathbf{J}_a$
  - integrability:  $\mathbf{F}^{ab} = 0$
- **Linking equations**
  - kinetic:  $\mathbf{A}^{ab} \xrightarrow{\chi} \mathbf{H}_a$
  - mass:  $\mathbf{A}^{ab} \xrightarrow{\xi} \mathbf{P}_{ab}$



# Analogy with massive electromagnetism

**Extrapolation** to Proca theory: impurities at the Planck scale  $M$  ?

$$L_{\text{em}} = \frac{1}{4} H^{\mu\nu} F_{\mu\nu} - \frac{1}{2} m^2 P^\mu A_\mu, \quad L_{\text{gr}} = \frac{1}{4} H_{\alpha\beta}^{\mu\nu} F^{\alpha\beta}_{\mu\nu} - \frac{1}{2} M^2 P_{\alpha\beta}^\mu A^{\alpha\beta}_\mu.$$

Objects and laws	Basis	Electromagnetism	Gravity	W	Electromagnetism	Gravity
Source current	3-	$\mathbf{J} = \mathbf{T} + \mathbf{t}$	$\mathbf{J}_a = \mathbf{T}_a + \mathbf{t}_a$	1	$J^\mu = T^\mu + t^\mu$	$J^\mu_\nu = T^\mu_\nu + t^\mu_\nu$
Conservation law	4-	$d\mathbf{J} = 0$	$d\mathbf{J}_a = 0$	1	$\nabla_\mu J^\mu = 0$	$\nabla_\mu J^\mu_\nu = 0$
Kinetic excitation	2-	$\mathbf{H}$	$\mathbf{H}_a, \mathbf{H}_{ab}$	1	$H^{\mu\nu}$	$H^{\mu\nu}_{\alpha\beta}, \mathbf{H}^{\mu\nu}_{\alpha\beta}$
Mass excitation	3-	$\mathbf{P}$	$\mathbf{P}_{ab}$	1	$P^\alpha$	$P^\alpha_{\mu\nu}$
Inhomog. field eqn.	3-	$d\mathbf{H} = \mathbf{J}$	$d\mathbf{H}_a = \mathbf{J}_a$	1	$\nabla_\mu H^{\mu\nu} = J^\nu$	$\nabla_\mu H^{\mu\nu}_\alpha = J^\nu_\alpha$
Kinetic potential	1+	$\mathbf{A}$	$\mathbf{A}^{ab}$	0	$A_\mu$	$A^\alpha_\nu$
Mass potential	0+	$\mathbf{B}$	$\mathbf{B}^{ab}$	0	$B$	$B^{\alpha\beta}$
Field strength	2+	$\mathbf{F} = d\mathbf{A}$	$\mathbf{F}^{ab} = d\mathbf{A}^{ab}$	0	$F_{\mu\nu} = 2\nabla_{[\mu} A_{\nu]}$	$F^{\alpha\beta}_{\mu\nu} = 2\nabla_{[\mu} A^{\alpha\beta}_{\nu]}$
Homog. field eqn	3+	$d\mathbf{F} = 0$	$d\mathbf{F}^{ab} = 0$	0	$\nabla_{[\alpha} F_{\mu\nu]} = 0$	$\nabla_{[\sigma} F^{\alpha\beta}_{\mu\nu]} = 0$
Lorentz force	4-	$f_a = \partial_a \cdot \mathbf{F} + \mathbf{J}$	$f_a^c = \partial_a \cdot \mathbf{F}^{cd} \wedge \mathbf{J}_d$	1	$f_\mu = F_{\mu\nu} J^\nu$	$f_\mu^\alpha = F^{\alpha\beta}_{\mu\nu} J^\nu_\beta$
Effective force	4-	$\mathbf{f} = -dt$	$\mathbf{f}_a = -dt_a$	1	$\mathbf{f} = -\nabla_\mu t^\mu$	$\mathbf{f}_\nu = -\nabla_\mu t^\mu_\nu$
Energy-momentum	3-	$\underset{\text{em}}{\mathbf{T}_a}$	$t_a$	1	$\underset{\text{em}}{\mathbf{T}_\mu^\nu}$	$t^\mu_\nu$
Kinetic Lagrangian	4-	$\overset{\text{kin}}{\Lambda} = -\frac{1}{2} \mathbf{F} \wedge \mathbf{H}$	$\overset{\text{kin}}{\Lambda} = -\frac{1}{2} \mathbf{F}^{ab} \wedge \mathbf{H}_{ab}$	1	$\overset{\text{kin}}{L} = \frac{1}{4} H^{\mu\nu} F_{\mu\nu}$	$\overset{\text{kin}}{L} = \frac{1}{4} H^{\mu\nu} \alpha_{\beta\gamma} F^{\alpha\beta}_{\mu\nu}$
Mass Lagrangian	4-	$\overset{\text{pot}}{\Lambda} = -\frac{1}{2} \mathbf{A} \wedge \mathbf{P}$	$\overset{\text{pot}}{\Lambda} = -\frac{1}{2} \mathbf{A}^{ab} \wedge \mathbf{P}_{ab}$	1	$\overset{\text{pot}}{L} = -\frac{1}{2} P^\mu A_\mu$	$\overset{\text{pot}}{L} = -\frac{1}{2} P^\alpha_{\mu\nu} A^{\mu\nu}_\alpha$
Kinetic constitutive rel.	0-	$\chi^{abcd}$	$\chi^{ab}_{\quad cdf} \chi^{ab}_{\quad cd\bar{f}}$	1	$\chi^{\alpha\beta\mu\nu}$	$\chi^{\mu\nu}_{\alpha\beta\gamma} \delta^{\delta}_{\alpha\beta\gamma\delta}$
Mass constitutive rel.	0-	$\xi^{ab}$	$\xi^a_{\quad cd} \xi^b_{\quad ef}$	1	$\xi^{\alpha\beta}$	$\xi^{\alpha}_{\mu\nu} \xi^{\beta}_{\mu\sigma}$

"The Standard Model is the result of the efforts to extend the ideas and methods of the electromagnetic interactions to all other forces in physics." -J. Iliopoulos, (2012).





# A new view at the Einstein equation

$$D\mathbf{H}_a = \mathbf{T}_a + \mathbf{t}_a$$

$\mathcal{G}^{\mu}_{\nu}$ Einstein tensor co-covariant	$\tau^{\mu}_{\nu}$ Gravitational EMT $\frac{1}{\sqrt{-g}} \nabla_{\alpha} (\sqrt{-g} H^{\alpha\mu\nu})$ $\rightarrow$ "Einstein EM complex"	$t^{\mu}_{\nu}$ Inertial EMT $\frac{1}{2} Q^{\mu}_{\nu} - P^{\mu}_{\alpha\beta} Q_{\nu}^{\alpha\beta}$ $\rightarrow \dot{t}^{\mu}_{\nu} = t^{\mu}_{E\nu}$	$T^{\mu}_{\nu}$ Matter EMT $-g_{\alpha\nu} \frac{2\delta(\sqrt{-g} L_M)}{\sqrt{-g} \delta g^{\mu\alpha}}$ co-covariant
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Jose Beltrán Jiménez, Lavinia Heisenberg, TSK: *The Canonical Frame of Purified Gravity* (IJMPD 2019):

The original  $H^{\mu\nu}_{\alpha}$  is vindicated in the now covariant formulation.  
The new conjecture: **the canonical frame** is  $t^{\mu}_{\nu} = 0$ .

- The field equations are both invariant and “pseudo-invariant”
- Their decomposition is not “pseudo-invariant”

# The energy and the entropy

Canonical frame:  $t_a = 0$

$$DH_a = T_a + \cancel{t_a}$$

*The physical principle that the metric should not contribute to the energy-momentum, determines uniquely the energetics and thermodynamics in a gravitating system.*

- **Energy-momentum** from Noether's theorem:

$$\mathring{P}_\mu = \int d^2S \sqrt{-g} \mathring{H}^{i0}{}_\mu n_i .$$

- **Entropy** from Euclidean QG:

$$-\log Z \approx \frac{1}{2} \int d^3V \sqrt{-g} \left( Q^\alpha - \tilde{Q}^\alpha \right) n_\alpha + \int d^4x \sqrt{-g} \left( L_M - \frac{1}{2} T \right) .$$

- **Consistency** from thermodynamics:

$$S = \beta E_0 + \log Z .$$

# Results for black holes (in natural units)

- The computation of Schwarzschild black hole entropy
  - in GR:  $S \approx \frac{1}{16\pi} \int d^4x \underbrace{\sqrt{-g}\mathcal{R}}_{=0} + \underbrace{\int d^3x \mathcal{K}}_{=\infty} - \underbrace{\int d^3x \mathcal{K}_0}_{=\infty} = 4\pi M^2.$
  - in Q-gravity:  $S \approx -\frac{1}{16\pi} \int d^3x \tilde{Q}^\mu n_\mu = 4\pi M^2.$
- For the energy we obtain consistently  $E_\mu = M\delta_\mu^0$ 
  - The computations are covariant, the results unique
  - ~~counter terms, boundary terms, higher derivatives, ambiguity~~
- Example 2: Reissner-Nordström black hole
  - The Euclidean action gives  $-\log Z \approx \pi r_+^2 + \beta q^2/(2r_+)$ ,
  - and we obtain for the canonical energy  $E_0 = M - q^2/(2r_+)$ .

# Cosmological results (in progress...)

- Example 3: de Sitter space
  - The Euclidean action gives  $-\log Z \approx \pi r_+^2$ ,
  - and we obtain the canonical energy-momentum  $E_\mu = 0$ .
- Example 4: general FLRW
  - I haven't found the solution in the canonical frame yet and
  - can only speculate about the Tryon hypothesis.
- Example 5: Gravitational waves
  - No  $\rho$  nor  $p$  but momentum:  $P_\mu = \frac{1}{2} \delta_\mu^j \int d^2S \sqrt{-g} h^i_j n_i$
  - Cooperstock hypothesis survives Feynman's bead argument.

# Hypercomplex and holographic magic



- The vacuum with vanishing energy density may have a finite energy charge.

- We couple to the vanishing  $\Gamma = \{ \} + L \stackrel{\text{def}}{=} 0$ , yet follow the geodesics of  $\{ \}$ .



- The integral of a vanishing action may describe a finite entropy.

# Conclusion

## Q-gravity is

 the canonical

- Geometrically:  $\mathring{Q} = -L_E$
- Materially: MCP ok

 translation gauge theory

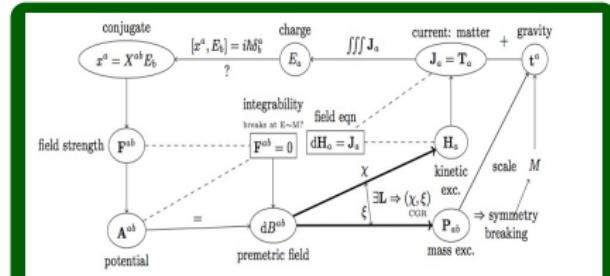
- TGR no more so than GR
- In STGR  $\Gamma$  is a translation

 which solves the  $E$ -problem

- Covariant energy
- Improved path integral

 and offers a new approach to resolve the conjugate  $t$ -problem.

- Complete analogy with Proca electromagnetism
- Generalisation of spacetime  $@\ell_{Planck}$ :  $Q^{ab} \rightarrow A^{ab} \neq -D\eta^{ab}$  !



$$L = -\frac{1}{2} F^{ab} \wedge H_{ab} - \frac{1}{2} M^2 A^{ab} \wedge P_{ab}$$